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A NONSTATIONARY AXISYMMETRIC MOTION OF GAS

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Axisymmetric nonstationary and irrotational motions of gas can be described by the system of equations [1],

$$\begin{aligned} r \frac{\partial a}{\partial t} + Nr \frac{\partial a}{\partial r} + T \frac{\partial a}{\partial \theta} + (\gamma - 1) a \left[\frac{\partial(Nr)}{\partial r} + \frac{\partial T}{\partial \theta} + N + T \operatorname{ctg} \theta \right] &= 0, \\ \frac{\partial N}{\partial t} + \frac{\partial}{\partial r} \left[(N^2 + T^2)/2 + a/(\gamma - 1) \right] &= 0, \\ \frac{\partial N}{\partial \theta} - \frac{\partial}{\partial r} (Tr) = 0, \quad a = \frac{dp}{d\rho}, \quad p = A\rho^\gamma, \end{aligned} \quad (1)$$

where N and T are the radial and the tangential components, respectively, of the gas velocity; a is the square of sound velocity; r, θ are the spherical coordinates.

A class of solutions will be found for system (1) assuming that the velocity components N, T depend on the angle θ and the time t only. It follows from the third equation in (1) that

$$N = f(\theta, t), \quad T = f'_\theta(\theta, t). \quad (2)$$

By inserting N and T as given by (2) into the second equation of (1), an expression is obtained for the square of sound velocity in terms of the function $f(\theta, t)$

$$a = -(\gamma - 1) [rf'_t + \psi(\theta, t)], \quad (3)$$

where $\psi(\theta, t)$ is an arbitrary function. The use of (2) and (3) reduces the first equation in (1) to the following system:

$$\begin{aligned} f = t\varphi(\theta) + x(\theta), \quad \psi + A(\theta)t^2/2 + B(\theta)t + \mu(\theta) &= 0, \\ (\gamma - 1)\psi(2f + f'_\theta \operatorname{ctg} \theta + f''_{\theta\theta}) + f'_\theta \psi'_\theta &= 0, \end{aligned} \quad (4)$$

where

$$A(\theta) = (2\gamma - 1)\varphi^2 + (\gamma - 1)\varphi''_{\theta\theta}\varphi + \varphi'^2 + (\gamma - 1)\varphi\varphi'_\theta \operatorname{ctg} \theta; \quad B(\theta) = (2\gamma - 1)x\varphi + (\gamma - 1)x''_{\theta\theta}\varphi + x'_\theta\varphi'_\theta + (\gamma - 1)\varphi x'_\theta \operatorname{ctg} \theta;$$

$\mu(\theta)$ is an arbitrary function. Since θ and t are independent variables, therefore (4) implies an overdetermined system of equations for finding $\varphi(\theta), x(\theta), \mu(\theta)$

$$\begin{aligned} (\gamma - 1)(2\varphi + \varphi'_\theta \operatorname{ctg} \theta + \varphi''_{\theta\theta})A + \varphi'_\theta A'_\theta &= 0, \\ (\gamma - 1)(2x + x'_\theta \operatorname{ctg} \theta + x''_{\theta\theta})A + x'_\theta A'_\theta + 2B(\gamma - 1)(2\varphi + \varphi'_\theta \operatorname{ctg} \theta + \varphi''_{\theta\theta}) + 2\varphi'_\theta B'_\theta &= 0, \\ (\gamma - 1)(2\varphi + \varphi'_\theta \operatorname{ctg} \theta + \varphi''_{\theta\theta})\mu + \varphi'_\theta \mu'_\theta + B(\gamma - 1)(2x + x'_\theta \operatorname{ctg} \theta + x''_{\theta\theta}) + x'_\theta B'_\theta &= 0, \\ (\gamma - 1)(2x + x'_\theta \operatorname{ctg} \theta + x''_{\theta\theta})\mu + x'_\theta \mu'_\theta &= 0. \end{aligned} \quad (5)$$

The consistency of (5) is now analyzed. By assuming that $\varphi'_\theta \neq 0, x'_\theta \neq 0$, one can eliminate from (5) $A'_\theta, B'_\theta, \mu'_\theta$. As a result, one arrives at a relation between x and φ ,

$$[Ax'_\theta/\varphi'_\theta + 2\mu\varphi'_\theta/x'_\theta - 2B][\Phi/\varphi'_\theta - X/x'_\theta] = 0, \quad (6)$$

where $\Phi = (2\varphi + \varphi'_\theta \operatorname{ctg} \theta + \varphi''_{\theta\theta})(\gamma - 1)$; $X = (2x + x'_\theta \operatorname{ctg} \theta + x''_{\theta\theta})(\gamma - 1)$.

By setting the second factor equal to zero, a differential relation is obtained between x and φ ,

$$x = \varphi \left[c_1 \int \varphi^{-2} R d\varphi + c_2 \right], \text{ where } R = \exp \left[2 \int \varphi (\varphi'_\theta)^{-1} d\theta \right]. \quad (7)$$

Under condition (7), system (5) is as follows:

for $A \neq 0$ one has

$$A\Phi + \varphi'_\theta A'_\theta = 0, \quad B = c_3 A, \quad \mu = c_4 A; \quad (8)$$

for $A = 0$ one has

$$(2\gamma - 1)\varphi^2 + (\gamma - 1)\varphi''_{\theta\theta}\varphi + \varphi'^2_{\theta} + (\gamma - 1)\varphi\varphi'_\theta \operatorname{ctg} \theta = 0, \quad (9)$$

$$\Phi B + \varphi'_\theta B'_\theta = 0, \quad \Phi\mu + \varphi'_\theta \mu'_\theta = 0.$$

The second equation of (8) imposes an additional differential relation on x and φ , which together with A and B , can be given by

$$A[c_3 - x'_\theta/\varphi'_\theta] - \varphi[x - \varphi x'_\theta/\varphi'_\theta] = 0. \quad (10)$$

By substituting the value of x from (7) into (10), one arrives at an equation for φ

$$A \left[\frac{c_3 - c_2}{c_1} - \int \varphi^{-2} R d\varphi - R/\varphi \right] - \varphi R = 0.$$

By differentiating the above with respect to θ , one obtains the first equation (8),

$$A\Phi + \varphi'_\theta A'_\theta = 0.$$

Thus, system (5) under condition (7) is consistent.

System (9) is now analyzed. The additional differential relation between x and φ in system (9) determines the second equation. The expression for B from (4) under the condition (7) and for $A = 0$ is

$$B = -c_1 \varphi R.$$

By substituting this value of B into the second equation in (9), one finds

$$R[(2\gamma - 1)\varphi^2 + (\gamma - 1)\varphi''_{\theta\theta}\varphi + \varphi'^2_{\theta} + (\gamma - 1)\varphi'_\theta \operatorname{ctg} \theta] = 0,$$

which implies that (9) is consistent. Then one has $\mu = c_3 \varphi R$ as well.

The case is now considered when the first factor in (6) vanishes, i.e.,

$$\mu = Bx'_\theta/\varphi'_\theta - (x'^2_{\theta} A)/(2\varphi'^2_{\theta}). \quad (11)$$

Inserting (11) in the last two equations of (5) one arrives at a system for x and φ

$$2B[X + F'_\theta x'_\theta/F] - A[FX + 2F'_\theta x'_\theta - \Phi F x'_\theta/\varphi'_\theta] + 2B'_\theta x'_\theta = 0, \quad (12)$$

$$B[\Phi F + \varphi'_\theta F'_\theta + X] - A\varphi'_\theta F F'_\theta + 2B'_\theta x'_\theta = 0,$$

where $F = x'_\theta/\varphi'_\theta$. By subtracting the first equation from the second, one finds that

$$[Ax'_\theta/\varphi'_\theta - B][x''_{\theta\theta}\varphi'_\theta - x'_\theta\varphi''_{\theta\theta}]/\varphi'_\theta + X - \Phi x'_\theta/\varphi'_\theta = 0. \quad (13)$$

By setting the first factor equal to zero, one finds

$$B = Ax'_\theta/\varphi'_\theta. \quad (14)$$

If one substitutes B from (14) into the second equation of (12), one arrives at the expression

$$A[F^2\Phi - FX + 2F'_\theta x'_\theta] = 0. \quad (15)$$

The equation $A = 0$ together with (14), (11) implies that $B = 0, \mu = 0$; the latter complies with the particular case of the condition (7). The vanishing of the second factor together with (15) yields an equation for x

$$(\gamma - 3)(x''_{\theta\theta}\varphi'_\theta - \varphi''_{\theta\theta}x'_\theta) + 2(\gamma - 1)(x\varphi'_\theta - \varphi x'_\theta) = 0,$$

the second equation for φ being given by (14). A consistent solution of these two equations is given by $x = c\varphi$, which also agrees with a particular case of condition (7).

If one sets the second factor in (13) equal to zero, one obtains a differential relation between x and φ ; employing the latter, the second equation in (12) is given in the form

$$(x'_\theta \varphi - x \varphi'_\theta)[(1 - 2\gamma)A + (3\gamma - 2)^2 \varphi^2 / \gamma] = 0.$$

The vanishing of the first factor implies that $x = c\varphi$. This relation between x and φ has previously been analyzed. The vanishing of the second factor results in an equation for φ ,

$$[2\gamma - 1 - (3\gamma - 2)^2 / (2\gamma^2 - \gamma)] \varphi^2 + \varphi_0^2 + (\gamma - 1) \varphi''_{\theta\theta} \varphi + (\gamma - 1) \varphi \varphi'_\theta \text{ctg } \theta = 0,$$

which together with the first equation (5) admits a consistent solution only given in the form $\varphi = 0$.

The case of $\varphi'_\theta = 0$, $x'_\theta = 0$ is now analyzed. For $\varphi'_\theta = 0$ Eq. (5) implies that $\varphi = 0$. By using (4) one obtains the stationary state of the flow,

$$N = x(\theta), T = x'_\theta(\theta), a^2 = (\gamma - 1)\mu(\theta).$$

Under the conditions of a stationary state, system (1) implies the Bernoulli-Euler integral which imposes a differential relation between x and μ given by $(x^2 + x'^2_\theta) / 2 + \mu = \text{const}$. If this equation is taken into account together with the last equation in (5), one can bring the system (1) to a single equation for the function $x(\theta)$,

$$x''_{\theta\theta} (a^2 - x'^2_\theta) + (2a^2 - x'^2_\theta) x + a^2 x'_\theta \text{ctg } \theta = 0, \\ a^2 = -(\gamma - 1)(x^2 + x'^2_\theta) / 2 + c,$$

which describes the axisymmetric conical gas flow; it was analyzed in [2, 3]. The case of $x'_\theta = 0$ ($x \neq 0$) results in $\mu = 0$, $B = 0$, $A = 0$. For $x = 0$ one has $B = 0$, $\mu = c_1 A$.

Thus, system (5), fully analyzed for consistency, is described in the case of a nonstationary state of gas flows by the following functions of the variables r , θ , t :

$$N(\theta, t) = \varphi \left[t + c_2 + c_1 \int \varphi^{-2} R d\varphi \right], \\ T(\theta, t) = \varphi'_\theta \left[t + c_2 + c_1 \left(\int \varphi^{-2} R d\varphi + \varphi^{-1} R \right) \right], \\ a(\theta, t) = -(\gamma - 1) \left[\varphi r + (c_3 t^2 / 2 + c_4 t + c_5) / (R \varphi'_\theta \sin \theta)^{(\gamma-1)} \right],$$

where φ can be found from the equation

$$\begin{aligned} & [(2\gamma - 1) \varphi^2 + (\gamma - 1) \varphi \varphi'_\theta \text{ctg } \theta + \varphi_0^2 + (\gamma - 1) \varphi \varphi''_{\theta\theta}] \Phi + \\ & + [(2\gamma - 1) \varphi^2 + (\gamma - 1) \varphi \varphi'_\theta \text{ctg } \theta + \varphi_0^2 + (\gamma - 1) \varphi \varphi''_{\theta\theta}]'_\theta \varphi'_\theta = 0. \end{aligned} \quad (16)$$

Two particular solutions of Eq. (16) are available. The first one is $\varphi = c_6 \cos(\theta + c_7)$. The other one can be found from the equation

$$(2\gamma - 1) \varphi^2 + (\gamma - 1) \varphi \varphi'_\theta \text{ctg } \theta + \varphi_0^2 + (\gamma - 1) \varphi \varphi''_{\theta\theta} = 0. \quad (17)$$

Equation (17) is transformed to a Legendre equation with the aid of the substitution $\varphi = y^{(\gamma-1)/\gamma}$. For the latter the general solution was found in [4] which for φ can be written as

$$\varphi = \left[c_6 F \left(-\frac{\nu}{2}, \frac{1+\nu}{2}, \frac{1}{2}, \cos^2 \theta \right) + c_7 \cos \theta \cdot F \left(\frac{1-\nu}{2}, 1 + \frac{\nu}{2}, \frac{3}{2}, \cos^2 \theta \right) \right]^{\frac{\gamma-1}{\gamma}},$$

where $\nu = (3\gamma - 2) / 2(\gamma - 1)$; F is the hypergeometric function.

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